# Halley's Method for the Matrix Sector Function * 

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#### Abstract

The matrix $n$-sector function is a generalization of the matrix sign function, and can be used to determine the number of eigenvalues of a matrix in a specific sector of the complex plane, and to extract the eigenpairs belonging to this sector without explicitly computing the eigenvalues. It is known that Newton's method, which can be used for computing the matrix sign function, is not globally convergent for the matrix sector function. The only existing algorithm for computing the matrix sector function is based on the continued fraction expansion approximation to the principal $n$th root of an arbitrary complex matrix. In this paper, we introduce a new algorithm, based on Halley's generalized iteration formula for solving nonlinear equations. It is shown that the iteration has good error propagation properties and high accuracy. Finally, we give two application examples, and summarize the results of our numerical experiments comparing Newton's, the continued fraction, and Halley's method.


## 1 Introduction

Fast computation of a restricted subset of eigenpairs of time varying matrices is an important topic in real time signal processing and control applications. Approximation of a matrix by another of lower rank, or model reduction, is desired in many applications, e.g., in systems theory [25], data analysis, pattern recognition, spectral analysis, radar, sonar, and geophysics [7]. Fast sequential and parallel for computing the eigenvalues and eigenvectors of unsymmetric matrices have been developed $[5,12,13,15]$. These algorithms can be used to compute all the eigenvalues of the matrix and then to extract the specified subset of the eigenvalue and eigenvector pairs. However, for these applications, methods which compute only a restricted subset of eigenpairs without resorting to computationally expensive eigenpair methods would be more useful and efficient. The matrix sign function and the matrix $n$-sector function can efficiently and reliably be used for this purpose [3, 2]. By obtaining the sector function of a matrix, we can easily determine the number of eigenvalues of a matrix in a specific sector of the complex plane, and extract the eigenpairs belonging to this sector without explicitly computing the eigenvalues and eigenvectors. By shifting the original matrix, or applying bilinear transformations, we can extend the sectors to various shapes and geometries. The matrix sign and sector function methods also have certain properties which make them more advantageous, e.g., complex arithmetic is avoided for matrices with real entries.

The matrix sign function is a particular case of the matrix $n$-sector function for $n=2$. Sequential and parallel algorithms for the matrix sign function have been developed [21, 16, 17, 11, 20], and its applications systems theory and matrix analysis have been established [23, 24, 6, 11]. However, the matrix sector function is a fairly new research topic. To the best of our knowledge, the only

[^0]existing algorithm for computing the sector function of arbitrary complex matrices is the one given in [27]. Let a matrix $A \in \mathcal{C}^{m \times m}$ have eigenspectrum $\sigma(A)=\left\{\lambda_{i}, i=1, \ldots, m\right\}$ where $\lambda_{i} \neq 0$ and $\arg \left(\lambda_{i}\right) \neq \pi$. The principal $n$th root of $A$, denoted as $\sqrt[n]{A} \in \mathcal{C}^{m \times m}$, is the matrix satisfying $(\sqrt[n]{A})^{n}=A$ and $\arg (\sigma(\sqrt[n]{A})) \in(-\pi / n, \pi / n)$. It is known that Newton's method can be used to compute the principal $n$th root of a positive definite matrix [14]. However, it has been pointed out in $[26,27]$ that Newton's method fails to give the principal $n$th root of a general complex matrix, thus, cannot be used for computing the matrix sector function.

In this paper, a fast and highly accurate algorithm is proposed. We first give the definitions of the matrix sector function and point out its applications. Then we introduce Halley's method, and give a perturbation analysis which shows that the first order errors in one step do not propagate to the next step. Finally, two application examples are given, and numerical experiments comparing Newton's, the continued fraction, and Halley's methods are summarized.

## 2 Definition of Matrix Sector Functions

The $n$-sector function of a scalar (matrix) is based on the principal $n$th root of the scalar (matrix). We begin with the definition of the sector function for a scalar. Let $\lambda \in \mathcal{C}$ be expressed by $\lambda=\rho e^{j \theta}$, where $\rho>0, j=\sqrt{-1}, \theta \in[0,2 \pi)$, and $\theta \neq 2 \pi\left(k+\frac{1}{2}\right) / n$ for $k \in[0, n-1]$. Assume that $\lambda$ lies within the sector $\Phi_{k}$ in $\mathcal{C}$ defined as the region bounded by the sector angles $2 \pi\left(k-\frac{1}{2}\right) / n$ and $2 \pi\left(k+\frac{1}{2}\right) / n$, where $k \in[0, n-1]$. Then the scalar $n$-sector function of $\lambda$ is defined as

$$
S_{n}(\lambda)=e^{j 2 \pi k / n}
$$

Let $\sqrt[n]{\lambda}$ be the principal $n$th root of $\lambda \in \mathcal{C}$. As shown in [27], the scalar sector function of $\lambda$ can be expressed as

$$
S_{n}(\lambda)=\frac{\lambda}{\sqrt[n]{\lambda^{n}}}
$$

where $\lambda \neq 0$ and $\arg (\lambda) \neq 2 \pi\left(k+\frac{1}{2}\right) / n$ for $k \in[0, n-1]$. Therefore, the scalar sector function maps a scalar in a specific sector to the bisector of the sector angles on the unit circle. The scalar sign function is a particular case of the $n$-sector function for $n=2$, i.e., the complex plane is partitioned into 2 sectors: Complex numbers with positive and negative real parts are mapped to +1 and -1 , respectively. Figure 1 shows the sector angles and the regions in the complex plane for $n=2,3,4$.


Figure 1: The sector regions and angles for $n=2,3,4$.
We also define the $q$ th $n$-sector function of the scalar $\lambda$ for $q \in[0, n-1]$, denoted by $S_{n, q}(\lambda)$, as the transformation that takes $\lambda$ to 1 if $\lambda$ belongs to $\Phi_{q}$ and to zero otherwise:

$$
S_{n, q}(\lambda)= \begin{cases}1 & \text { if } \lambda \in \Phi_{q} \\ 0 & \text { otherwise }\end{cases}
$$

We can extend these definitions to complex square matrices as follows. Let $A \in \mathcal{C}^{m \times m}$ and $\sigma(A)=$ $\left\{\lambda_{i}, i=1, \ldots, m\right\}$ be its spectrum with not necessarily distinct eigenvalues $\lambda_{i} \neq 0$ and $\arg \left(\lambda_{i}\right) \neq$ $2 \pi\left(k+\frac{1}{2}\right) / n$ for $k \in[0, n-1]$. Let $M \in \mathcal{C}^{m \times m}$ be the modal matrix that takes $A$ to its Jordan form as

$$
A=M\left[J_{1} \oplus J_{2} \oplus \cdots \oplus J_{k}\right] M^{-1},
$$

where $J_{i} \in \mathcal{C}^{r_{i} \times r_{i}}$ are the Jordan blocks corresponding to the $i$ th eigenvalue with geometric multiplicity $r_{i}$, such that $\sum_{i=1}^{k} r_{i}=m$. Applying the matrix function definition of Giorgi [22], we can define the matrix sector function of $A$ as

$$
\begin{aligned}
S_{n}(A) & =A\left(\sqrt[n]{A^{n}}\right)^{-1} \\
& =M\left[J_{1}\left(\sqrt[n]{J_{1}^{n}}\right)^{-1} \oplus \cdots \oplus J_{k}\left(\sqrt[n]{J_{k}^{n}}\right)^{-1}\right] M^{-1}
\end{aligned}
$$

where

$$
J_{i}\left(\sqrt[n]{J_{i}^{n}}\right)^{-1}=S_{n}\left(\lambda_{i}\right) I_{r_{i}}
$$

Therefore, the definition of the sector function of a matrix becomes

$$
S_{n}(A)=M\left[\bigoplus_{i=1}^{m} S_{n}\left(\lambda_{i}\right)\right] M^{-1}
$$

Following the definition of the scalar sector function, we see that the matrix sector function maps the eigenvalues of a given matrix to the bisector of the sector angles of the corresponding region onto the unit circle while preserving the eigenvectors. Similarly, the matrix $q$ th $n$-sector function of $A$, denoted by $S_{n, q}(A)$, is defined as

$$
S_{n, q}(A)=M\left[\bigoplus_{i=1}^{m} S_{n, q}\left(\lambda_{i}\right)\right] M^{-1}
$$

where $S_{n, q}(\lambda)$ is the scalar $q$ th $n$-sector function of $\lambda$. The matrix $q$ th $n$-sector function of $A$ maps the eigenvalues of $A$ in the sector $\Phi_{q}$ to 1 , and the remaining eigenvalues are mapped to zero. It can be easily proven (see, Theorem 4.2 in [26]) that the matrix $q$ th $n$-sector function of $A$ is equal to

$$
S_{n, q}(A)=\frac{1}{n} \sum_{i=1}^{n}\left[S_{n}(A) e^{-j 2 \pi q / n}\right]^{i-1}
$$

for $q \in[0, n-1]$, where $S_{n}(A)$ is the $n$-sector function of $A$.

## 3 Applications of Matrix Sector Functions

The matrix sector function can be utilized to block diagonalize a given matrix, without explicitly computing the eigenvalues and the corresponding eigenvectors. For $A \in \mathcal{C}^{m \times p}$, we define $\operatorname{ind}[A]$ as the set of linearly independent column vectors of $A$. Let $\mu_{i}$ denote the linearly independent column vectors of $S_{n, q}(A)$ for $A \in \mathcal{C}^{m \times m}$ with nonzero eigenvalues, i.e.,

$$
\mu_{i}=\operatorname{ind}\left[S_{n, q}(A)\right] \in \mathcal{C}^{m \times m_{i}}
$$

for $q \in[0, n-1], i \in[1, k]$, and $m=\sum_{i=1}^{k} m_{i}$. The block modal matrix $M$, defined as

$$
M=\left[\mu_{1}, \mu_{2}, \ldots, \mu_{k}\right] \in \mathcal{C}^{m \times m}
$$

can be used to block diagonalize the matrix $A$ as

$$
D=M^{-1} A M=\operatorname{diag}\left[A_{1}, A_{2}, \ldots, A_{k}\right]
$$

where the block elements correspond to the eigenvalues in the specified sector of the complex plane. This strategy can be used to decompose a system into several smaller subsystems with similar transient characteristics. The location of the poles with respect to the sector angles determines the natural frequencies and the damping ratio of the system [1]. Decoupling with respect to the given sectors would enable us to obtain a physical realization which is more precise and stable. This analysis can be performed from both state-space and matrix-fraction description points of view. Let a $q$-input, $p$-output system be described by

$$
\begin{aligned}
& \dot{x}(t)=A x(t)+B u(t) \\
& y(t)=C x(t)+D u(t)
\end{aligned}
$$

where $x(t) \in \mathcal{C}^{m \times 1}, u(t) \in \mathcal{C}^{q \times 1}$, and $y(t) \in \mathcal{C}^{p \times 1}$. Assuming the system is observable and controllable, we can define the left and right matrix fraction description of the system as

$$
\begin{aligned}
& H_{l}(s)=C(s I-A)^{-1} B+D=D_{l}^{-1}(s) N_{l}(s)+D \\
& H_{r}(s)=C(s I-A)^{-1} B+D=N_{r}^{-1}(s) D_{r}(s)+D
\end{aligned}
$$

where $N_{r}(s), D_{r}(s), N_{l}(s)$, and $D_{l}(s)$ are polynomial matrices. Let $M$ be the block modal matrix which block diagonalizes $A$, obtained using the matrix sector functions $S_{n, q}(A)$. We have

$$
\begin{aligned}
A_{d} & =M^{-1} A M=\operatorname{diag}\left[\hat{A}_{1}, \hat{A}_{2}, \ldots, \hat{A}_{k}\right] \text { for } \hat{A}_{i} \in \mathcal{C}^{m_{i} \times m_{i}} \\
B_{d} & =M^{-1} B=\left[\hat{B}_{1}^{T}, \hat{B}_{2}^{T}, \ldots, \hat{B}_{k}^{T}\right]^{T} \text { for } \hat{B}_{i} \in \mathcal{C}^{m_{i} \times q} \\
C_{d} & =C M=\left[\hat{C}_{1}, \hat{C}_{2}, \ldots, \hat{C}_{k}\right] \text { for } \hat{C}_{i} \in \mathcal{C}^{p \times m_{i}}
\end{aligned}
$$

Thus, the system can be block decomposed into $k$ subsystems as

$$
\begin{aligned}
\dot{x}_{d}(t) & =A_{d} x_{d}(t)+B_{d} u(t) \\
y(t) & =C_{d} x_{d}(t)+D u(t)
\end{aligned}
$$

where $x(t)=M x_{d}(t)$. The input/output relationship after the decomposition is given as

$$
Y(s)=\left(C_{d}\left(s I-A_{d}\right)^{-1} B_{d}+D\right) U(s)
$$

where the $i$ th element of the transfer function matrix contains the $m_{i}$ eigenvalues of the sector $\Phi_{i}$.

## 4 Halley's Method for the Matrix Sector Function

It has been shown that Halley's generalized iteration formula for solving nonlinear equations is of third order, and its error-cubing variation converges faster than Newton's method [10, 9, 8, 4]. Halley's method can be derived by applying Newton's method to the function

$$
g(s)=\frac{f(s)}{\sqrt{f^{\prime}(s)}}
$$

which is written as

$$
\begin{equation*}
s_{k+1}=s_{k}-\frac{f\left(s_{k}\right)}{f^{\prime}\left(s_{k}\right)-\frac{f^{\prime \prime}\left(s_{k}\right) f\left(s_{k}\right)}{2 f^{\prime}\left(s_{k}\right)}} . \tag{1}
\end{equation*}
$$

We start with an alternative definition of the scalar sector function as the solution of the following equation

$$
f(s)=s^{n}-1=0 .
$$

Solution of this equation with Halley iteration becomes

$$
s_{k+1}=s_{k}-\frac{2 s_{k}\left(s_{k}^{n}-1\right)}{(n+1) s_{k}^{n}+(n-1)},
$$

which reduces to

$$
s_{k+1}=s_{k} \frac{(n-1) s_{k}^{n}+(n+1)}{(n+1) s_{k}^{n}+(n-1)} .
$$

This iteration produces an order $[1,1]$ rational Padé approximant to

$$
f(s)=\frac{s}{\sqrt[n]{1-z}}
$$

where $z=1-s^{n}$. In the matrix case, $f\left(S_{n}\right)$ can be defined as

$$
f\left(S_{n}\right)=S_{n}^{n}(A)-I=0 .
$$

Let $S_{n}[k]$ stand for the value of $S_{n}(A)$ at step $k$. Applying the iteration of Equation (1) to the above expression, we obtain Halley's method for matrix sector function as

$$
\begin{align*}
S_{n}[0] & =A, \\
S_{n}[k+1] & =S_{n}[k] \times\left((n-1) S_{n}^{n}[k]+(n+1) I\right) \times\left((n+1) S_{n}^{n}[k]-(n-1) I\right)^{-1},  \tag{2}\\
\lim _{k \rightarrow \infty} S_{n}[k] & =S_{n}(A) .
\end{align*}
$$

Now we give a convergence analysis of Halley's method by checking the location of the eigenvalues of $S_{n}(A)$ as the algorithm iterates starting from $k=0$, i.e., $S_{n}[0]=A$. We assume that $S_{n}[k]$ has an eigenvalue $\lambda_{k}$ in the sector $\Phi_{q}$ at the $k$ th step of the iteration, which can be expressed as

$$
\lambda_{k}=\rho_{k} e^{j\left(\psi_{k}+2 \pi q / n\right)},
$$

where $\rho_{k}=\left|\lambda_{k}\right|, q \in[0, n-1]$, and $\left|\psi_{k}\right|<\pi / n$. Here, $\lambda_{k+1}$ can be given as

$$
\lambda_{k+1}=\rho_{k} e^{j\left(\psi_{k}+2 \pi q / n\right)} \frac{(n-1) \rho_{k}^{n} e^{j n \psi_{k}}+(n+1)}{(n+1) \rho_{k}^{n} e^{j n \psi_{k}}+(n-1)} .
$$

We expect

$$
\lim _{k \rightarrow \infty} \rho_{k}=1 \text { and } \lim _{k \rightarrow \infty} \psi_{k}=0 .
$$

Let $\lim _{k \rightarrow \infty} \lambda_{k}$ exist and be finite. Denoting this limit by $x$, from Equation (2) we obtain

$$
x=x \frac{(n-1) x^{n}+(n+1)}{(n+1) x^{n}+(n-1)}
$$

which reduces to

$$
x^{n+1}-x=0 .
$$

Assuming $x \neq 0$, we find the solution of the limit equation as $x=\sqrt[n]{1}$, i.e., an $n$th root of unity.
Let $z$ be a complex number in the sector $\Phi_{0}$ which contains the first real root of unity, +1 . For the $n$ sector plane the following inequality should hold:

$$
\begin{equation*}
\left|\frac{(z-1)}{\left(z-z_{i}\right)}\right|<1 \tag{3}
\end{equation*}
$$

for $1 \leq i \leq n-1$, where $z_{i}=e^{j(2 \pi i) / n}$, i.e., one of the $n$th roots of unity. In order to guarantee that the consecutive iterates does not pass the sector boundaries, this inequality should hold true at each step of the iteration. Let $s_{k}$ be a scalar at the $k$ th step of the iteration, then the following equality should also be satisfied for all $k$ [19].

$$
s_{k+1}-S_{n}\left(s_{k+1}\right)=s_{k+1}-S_{n}\left(s_{k}\right) .
$$

The relationship between two consecutive iterates is obtained as

$$
\frac{\left(s_{k+1}-1\right)}{\left(s_{k+1}-z\right)}=\frac{\left(s_{k}-1\right)^{3}}{\left(s_{k}-z\right)^{3}}\left(\frac{\sum_{j=1}^{n-1}\left(\sum_{i=1}^{j}(n-(2 i-1)) s_{k}^{n-j-1}\right)}{\sum_{j=1}^{n-1}\left(\sum_{i=1}^{j}(n-(2 i-1)) s_{k}^{n-j-1} z^{j-1}\right)}\right)
$$

where $z$ is one of the $n$th roots of unity. The first part of right hand side satisfies the convergence property (3), but the entire right hand side contains a rational term which may force the ratio to be greater than 1 for some $s_{k}$. In this case the iterate changes sectors and the iteration converges to an incorrect value. The following example was provided by Kenney and Laub [18]: Consider $z_{1}=$ $e^{2 \pi / 3 j}$, i.e., the second cuberoot of unity. Taking $s_{0}=-z_{1}+0.001=0.5010-0.8660 j$ (which is in the sector $\Phi_{0}$ ) as our initial point the first step of the Halley iteration yields $s_{1}=-0.4920+0.8660 j$, and eventually the iteration converges to $z_{1}$ rather than to 1 . Such inaccuracies mostly occur for points on or near the sector boundaries.

## 5 Perturbation Analysis

In this section, we analyze Halley's iteration for matrix sector functions when the iterates are subject to perturbations from rounding errors at a given step $k$. Let $\tilde{S}_{n}[k]=S_{n}[k]+E[k]$, where $E[k]$ is the error at step k . The perturbed value of $S_{n}[k+1]$ can be written as

$$
\begin{aligned}
\tilde{S}_{n}[k+1]= & \left(S_{n}[k]+E[k]\right) \times\left((n-1)\left(S_{n}^{n}[k]+F[k]\right)+(n+1) I\right) \times \\
& \left((n+1)\left(S_{n}^{n}[k]+F[k]\right)+(n-1) I\right)^{-1},
\end{aligned}
$$

where

$$
F[k]=E[k] S_{n}^{n-1}[k]+S_{n}[k] E[k] S_{n}^{n-2}[k]+\cdots+S_{n}^{n-1} E[k] .
$$

Here, we have used the power expansion

$$
(A+E)^{n} \cong A^{n}+E A^{n-1}+A E A^{n-2}+A^{2} E A^{n-3}+\cdots+A^{n-1} E
$$

by ignoring the terms involving more than one error term. Assuming

$$
\left\|(n+1) S_{n}^{n}[k]+(n-1) I\right\|>\|(n+1) F[k]\|,
$$

we utilize the perturbation formula in [28] and obtain

$$
(A+E)^{-1}=A^{-1}-A^{-1} E A^{-1}+O\left(\|E\|^{2}\right) .
$$

Let

$$
\begin{aligned}
N[k] & =(n-1) S_{n}^{n}[k]+(n+1) I, \\
D[k] & =(n+1) S_{n}^{n}[k]+(n-1) I .
\end{aligned}
$$

Ignoring error terms of degree two or more, we obtain

$$
\begin{aligned}
\tilde{S}_{n}[k+1]= & S_{n}[k] N[k] D^{-1}[k]-(n+1) S_{n}[k] N[k] D^{-1}[k] F[k] D^{-1}[k]+ \\
& (n-1) S_{n}[k] F[k] D^{-1}[k]+E[k] N[k] D^{-1}[k] .
\end{aligned}
$$

This gives the error expression as

$$
\begin{aligned}
E[k+1] & =\tilde{S}_{n}[k+1]-S_{n}[k+1] \\
& =E[k] N[k] D^{-1}[k]+(n-1) S_{n}[k] F[k] D^{-1}[k]-(n+1) S_{n}[k+1] F[k] D^{-1}[k]
\end{aligned}
$$

Let $M$ be the modal matrix of $S_{n}[k]$ and $S_{n}[k+1]$ such that

$$
\begin{aligned}
D[k] & =M^{-1} S_{n}[k] M
\end{aligned}=\operatorname{diag}\left(\lambda_{1}[k], \ldots, \lambda_{m}[k]\right),
$$

Let $\hat{E}[k]=M^{-1} E[k] M$ and $\hat{F}[k]=M^{-1} F[k] M$, then $\hat{F}[k]$ can be written elementwise as

$$
\begin{equation*}
\hat{F}_{i j}[k]=\left(\sum_{l=0}^{n-1} \lambda_{i}^{l}[k] \lambda_{j}^{n-1-l}[k]\right) \hat{E}_{i j}[k] \tag{4}
\end{equation*}
$$

The matrix sector function identity suggests that

$$
\lim _{k \rightarrow \infty} S_{n}^{n}[k]=I
$$

thus, the elementwise error iteration becomes

$$
\hat{E}_{i j}[k+1]=\hat{E}_{i j}[k]+\frac{(n-1)}{2 n} \lambda_{i}[k] \hat{F}_{i j}[k]-\frac{(n+1)}{2 n} \lambda_{i}[k+1] \hat{F}_{i j}[k]
$$

where $\lim _{k \rightarrow \infty} \lambda_{i}[k]=e^{j 2 \pi q / n}$. Therefore, the error expression becomes

$$
\begin{equation*}
\hat{E}_{i j}[k+1]=\hat{E}_{i j}[k]-\frac{1}{n} e^{j 2 \pi q / n} \hat{F}_{i j}[k] \tag{5}
\end{equation*}
$$

We have two cases to consider:

1. If $\lambda_{i}$ and $\lambda_{j}$ are in different sectors, Equation (4) gives $\hat{F}_{i j}[k]=0$, and the error expression (5) becomes

$$
\hat{E}_{i j}[k+1]=\hat{E}_{i j}[k]
$$

i.e., the error in the $k$ th step is passed to, but not magnified in the $(k+1)$ st step.
2. If $\lambda_{i}$ and $\lambda_{j}$ are in the same sector, Equation (4) gives

$$
\hat{F}_{i j}[k]=n e^{j 2 \pi q(n-1) / n} \hat{E}_{i j}[k]
$$

Thus, the error expression (5) becomes

$$
\hat{E}_{i j}[k+1]=0
$$

i.e., no first order error is propagated to the $(k+1)$ st step.

Thus, we conclude that the first order errors in one step either have limited effect on the next step, or do not propagate to the next step at all.

## 6 Numerical Experiments

In this section, we analyze the accuracy of the three algorithms, namely, Newton's method, the continued fraction method, and Halley's method, according to the location of the eigenvalues with respect to the sector angles. We start with a diagonal matrix

$$
D=\operatorname{diag}(7+7 j, 7-7 j, 70+70 j, 70-70 j),
$$

whose eigenvalues are on the sector angles $\mp \pi / 4$ of the 4 -sector plane. We shift the real part of this matrix by $1 / \mu$ for $\mu>0$ to obtain $D(\mu)$ as

$$
D(\mu)=\operatorname{diag}\left(\left(7+\frac{1}{\mu}\right)+7 j,\left(7+\frac{1}{\mu}\right)-7 j,\left(70+\frac{1}{\mu}\right)+70 j,\left(70+\frac{1}{\mu}\right)-70 j\right) .
$$

This changes the location of the eigenvalues of this matrix from the sector angles to $\Phi_{0}$ on the 4 -sector plane. After this small shift, we form the following upper triangular matrix

$$
A(\mu)=D(\mu)+T
$$

by adding a strictly upper triangular matrix $T$, with elements uniformly distributed over the interval $[0,1]$. Now, an accurate sector function algorithm should produce $S_{4}(A(\mu))$ with all eigenvalues equal to 1 , i.e.,

$$
\sigma\left(S_{4}(A(\mu))\right)=\{1,1,1,1\} \text { for all } \mu>0 .
$$

We have applied Newton's, the continued fraction, and Halley's methods to compute $S_{4}(A(\mu))$ for several values of $\mu$. The results are summarized in Table 1.

| $\mu$ | Newton |  | Continued Fraction |  | Halley |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $k$ | $\sigma\left(S_{4}[k]\right)$ | $k$ | $\sigma\left(S_{4}[k]\right)$ | $k$ | $\sigma\left(S_{4}[k]\right)$ |
| 1 | 35 | $\pm j$ | 20 | 1 | 17 | 1 |
| 10 | 47 | $\pm 1$ | 22 | $1, \pm j$ | 20 | 1 |
| $10^{2}$ | 49 | $1, \pm j$ | 26 | $\pm j$ | 23 | 1 |
| $10^{3}$ | 54 | $\pm j$ | 29 | $1, \pm j$ | 25 | 1 |
| $10^{4}$ | 65 | $\pm j$ | 32 | $\pm j$ | 28 | 1 |
| $10^{5}$ | 71 | 1 | 35 | $\pm j$ | 31 | 1 |

Table 1. The eigenvalues of $S_{4}(A)$ after the convergence.
Our experiments shows that only Halley's iteration converges accurately for all $\mu$. Furthermore, Halley's method computes $S_{4}(A)$ using fewer iterations than both Newton's and the continued fraction methods. We also note that the iterative matrix sector algorithm may use more floatingpoint operations than the QR algorithm, but it is easier to parallelize and contains simple matrix operations such as LU decomposition and matrix multiplication.

## 7 An Application Example

Here we show how the impulse response of a system can be decomposed into its oscillatory and damped exponential components by utilizing the matrix sector functions. We consider a linear time
invariant system, represented by the matrices

$$
A=\left[\begin{array}{rrrrr}
-2.7798 & 14.4361 & -11.9801 & -28.2392 & 27.3195 \\
-5.1596 & 28.2055 & -15.6936 & -52.2117 & 46.9724 \\
3.4870 & -5.8394 & -9.6580 & 7.5337 & 0.1969 \\
7.2000 & -20.0000 & -2.2000 & 27.8000 & -20.0000 \\
12.3290 & -41.4465 & 3.2807 & 65.0112 & -50.2677
\end{array}\right],
$$

and

$$
B=\left[\begin{array}{lllll}
12 & 6 & 12 & 13 & 16
\end{array}\right]^{T}, \quad C=\left[\begin{array}{llllll}
0.1710 & -4.7202 & 4.8860 & 9.1554 & -9.3990
\end{array}\right] .
$$

The eigenvalues of the open-loop system are $\sigma(A)=\{-1 \pm 3.87 j,-1.5,-1.6 \pm 1.2 j\}$, therefore, the impulse response has a damped exponential and an oscillatory component. Computing $S_{4}(A)$ enables us to decompose the system into two components $A_{1} \in \mathcal{C}^{2 \times 2}$ and $A_{2} \in \mathcal{C}^{3 \times 3}$, with eigenvalues $\sigma\left(A_{1}\right)=\{-1 \pm 3.87 j\}$ and $\sigma\left(A_{2}\right)=\{-1.5,-1.6 \pm 1.2 j\}$. On the 4 -sector plane, the eigenvalues of $A_{1}$ lie in $\Phi_{1}$ and $\Phi_{3}$, where the damping ratio $\xi$ is less than 0.707 , and the eigenvalues of $A_{2}$ lies in $\Phi_{2}$, where $\xi>0.707$. Figure 2 shows the total and decomposed system impulse responses, in which the solid and dashed lines correspond to the impulse response of subsystems $A_{1}$ and $A_{2}$, respectively.


Figure 2: The total and decomposed system impulse responses.

## 8 Conclusion

We have described an iterative algorithm for the computation of the matrix sector function, which is based on the solution of a nonlinear equation using Halley's method. The algorithm is fast and numerically stable, and gives accurate results even for matrices with ill-conditioned eigenstructures. We have discussed applications of matrix sector functions, and provided some examples supporting these applications. We are currently investigating the effects of scaling on the speed of convergence, and developing efficient methods for the computation of partitioned matrix sector functions from the matrix sector function.

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## References

[1] A. A. Abdul-Wahab. Lyapunov-type equations for matrix root-clustering in subregions of the complex plane. International Journal of Systems Science, 21(9):1819-1830, 1990.
[2] Z. Bai and J. Demmel. Design of parallel nonsymmetric eigenvalue toolbox, Part 1. Technical Report csd-92-718, University of California Berkeley, December 1992.
[3] Z. Bai, J. Demmel, and M. Gu. Inverse free parallel spectral divide and conquer algorithms for nonsymmetric eigenproblems. Technical Report csd-94-793, University of California Berkeley, February 1994.
[4] S. Bittanti, A. J. Laub, and J. C. Willems, editors. The Riccati Equation. New York, NY: Springer-Verlag, 1991.
[5] D. Boley. Solving the generalized eigenvalue problem on a synchronous linear processor array. Parallel Computing, 3(2):153-166, May 1986.
[6] R. Byers. Solving the algebraic Riccati equation with the matrix sign function. Linear Algebra and its Applications, 85:267-279, 1987.
[7] P. Comon. Fast computation of a restricted subset of eigenpairs of a varying Hermitian matrix. In G. H. Golub and P. van Dooren, editors, Numerical Linear Algebra, Digital Signal Processing and Parallel Algorithms, pages 457-466. New York, NY: Springer-Verlag, 1991.
[8] A. A. M. Cuyt. Numerical stability of the Halley-iteration for the solution of a system of nonlinear equations. Mathematics of Computation, 38(157):171-179, January 1982.
[9] M. Davies and B. Dawson. On the global convergence of Halley's iteration formula. Numerische Mathematik, 14(2):133-135, 1975.
[10] J. S. Frame. A variation of Newton's method. The American Mathematical Monthly, 51:36-38, 1944.
[11] J. D. Gardiner and A. J. Laub. Parallel algorithms for algebraic Riccati equations. International Journal of Control, 44(6):1317-1333, December 1991.
[12] G. A. Geist, R. C. Ward, G. J. Davis, and R. E. Funderlic. Finding eigenvalues and eigenvectors of unsymmetric matrices using a hypercube multiprocessor. In G. Fox, editor, Hypercube Concurrent Computers and Applications, volume II, pages 1577-1582. New York, NY: ACM Press, 1988.
[13] G. H. Golub and C. F. van Loan. Matrix Computations. Baltimore, MD: The Johns Hopkins University Press, 2nd edition, 1989.
[14] W. D. Hoskins and D. J. Walton. A faster, more stable method for computing the $p$ th roots of positive definite matrices. Linear Algebra and its Applications, 26:139-163, 1979.
[15] A. Jennings and J. J. McKeown. Matrix Computations. New York, NY: John Wiley \& Sons, 1992.
[16] C. Kenney and A. J. Laub. Polar decomposition and matrix sign function condition estimates. SIAM Journal on Scientific and Statistical Computing, 12(3):488-504, 1991.
[17] C. Kenney and A. J. Laub. Rational iterative methods for the matrix sign function. SIAM Journal on Matrix Analysis and Applications, 12(2):273-291, April 1991.
[18] C. Kenney and A. J. Laub. Private communication, 1994.
[19] C. Kenney and A. J. Laub. A hyperbolic tangent identity and geometry of Padé sign function iterations. Numerical Algorithms, 7:111-128, 1994.
[20] Ç. K. Koç, B. Bakkaloğlu, and L. S. Shieh. Computation of the matrix sign function using continued fraction expansion. IEEE Transactions on Automatic Control, 39(8):1644-1647, August 1994.
[21] P. Pandey, C. Kenney, and A. J. Laub. A parallel algorithm for the matrix sign function. International Journal of High-Speed Computing, 2(2):181-191, 1990.
[22] R. F. Rinehart. The equivalence of definitions of a matric function. The American Mathematical Monthly, 3(62):395-414, 1955.
[23] J. D. Roberts. Linear model reduction and solution of the algebraic Riccati equation by use of the sign function. International Journal of Control, 32(4):677-687, 1980.
[24] L. S. Shieh, H. M. Dib, and R. E. Yates. Separation of matrix eigenvalues and structural decomposition of large-scale systems. IEE Proceedings: Control Theory and Applications, 133(2):90-96, 1986.
[25] L. S. Shieh, J. S. H. Tsai, and R. E. Yates. The generalized matrix sector functions and their applications to systems theory. IMA Journal of Mathematical Control and Information, 2:251-258, 1990.
[26] L. S. Shieh, Y. T. Tsay, and C. T. Wang. Matrix sector functions and their applications to systems theory. IEE Proceedings: Control Theory and Applications, 131(5):171-181, September 1984.
[27] J. S. H. Tsai, L. S. Shieh, and R. E. Yates. Fast and stable algorithms for computing the principal $n$th root of a complex matrix and the matrix sector functions. Computers and Mathematics with Applications, 15(11):903-913, 1988.
[28] J. H. Wilkinson. The Algebraic Eigenvalue Problem. Oxford, London: Oxford University Press, 1965.


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